

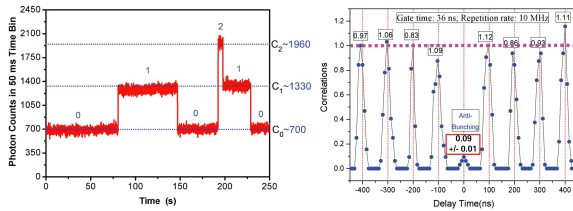
# 852-nm Single Photon Sources Based on Single Cesium Atom and Cesium Rydberg Super-atom Trapped in Optical Tweezer

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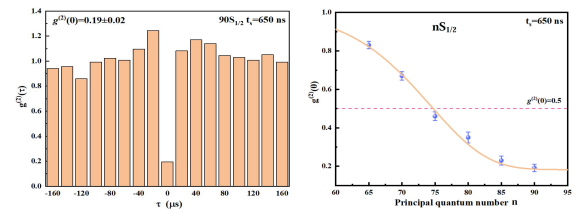
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Narrow-band 852-nm single-photon sources (SPS) are demonstrated based on single cold Cesium (Cs) atom and cold Cs Rydberg super-atom trapped in 1064-nm optical tweezer. For single atom approach, a single Cs atom is loaded into a 1064-nm far-off-resonance microscopic tweezer (FORT) from a single-atom magneto-optical trap (MOT), and prepared at the  $6S_{1/2} |F=4, m=+4\rangle$  state and subsequently has been periodically excited by the high-repetition resonant 852-nm 5-ns pulse chain to the  $6P_{3/2} |F'=5, m=+5\rangle$  state. Spontaneously radiated photons from  $|5', +5\rangle$  to  $|4, +4\rangle$  transition are collected and characterized by Hanbury Brown - Twiss (HBT) measurement. Strong photon antibunching,  $g^{(2)}(\tau=0) = 0.09 \pm 0.01$ , is demonstrated, and it confirms the single-photon characters. For Cs  $|5', +5\rangle$  to  $|4, +4\rangle$  cycling transition at 852 nm, the FORT magic wavelength has been calculated theoretically and verified experimentally to be 937.6 nm. The Hong-Ou-Mandel (HOM) two-photon interference measurement is utilized to evaluate the photon indistinguishability of our two SPS with 1064-nm FORT or 937.6-nm FORT. Preliminary experimental results in the 937.6-nm magic-wavelength FORT case indicate that the photon indistinguishability can be improved  $\sim 20\%$ , compared with the 1064-nm FORT case.

For Rydberg super-atom approach,  $\sim 10^4$  cold Cs atoms are loaded into a tightly-focused ellipse 1064-nm FORT from cold atomic ensemble after MOT + compress MOT + molasses phases, and prepared at the  $6S_{1/2} |F=4, m=0\rangle$  ground state. During the period for FORT switching off,  $\sim 100$  Cs atoms in a transverse region with  $\sim 10 \mu\text{m}$  diameter inside the ellipse cold atomic cloud are excited by means of tightly-focused two-photon resonant 852-nm and 509-nm laser pulses with a single-photon detuning of  $+30\text{MHz}$ . In the case of the principal quantum number  $n$  of  $nS_{1/2}$  Rydberg state is bigger enough, for example,  $90S_{1/2}$ , the Rydberg excitation blockade radius is larger than the 852-nm and 509-nm beam spots size ( $\sim 10 \mu\text{m}$ ), only one of  $\sim 100$  Cs atoms can be populated at  $90S_{1/2}$  Rydberg state because of Rydberg excitation blockade, and other atoms remain at  $6S_{1/2} |4, 0\rangle$  ground state. This is so called the single collective excitation Rydberg super-atom. After 650-ns storage time, the resonant 509-nm laser pulse reads out 852-nm single photon from Rydberg super-atom. The second-order correlation degree is measured to be  $g^{(2)}(\tau=0) = 0.19 \pm 0.02$  by the HBT scheme, and this confirms single-photon characters. Also magic-wavelength tweezer for Cs  $6S_{1/2} |4, 0\rangle$  ground state and  $90S_{1/2}$  Rydberg state is calculated theoretically. The magic-wavelength FORT for Rydberg super-atom approach will be implemented in the near future, and other efforts will be made to improve the photon purity, photon indistinguishability, and the photon extraction efficiency of SPS.



**Fig. 1:** Single-atom tweezer case. Left: the fluorescence photon counts of Cesium MOT. Right: the normalized correlations vs the delay time.  $g^{(2)}(\tau=0)=0.09\pm0.01$ .



**Fig. 2:** Cesium Rydberg super-atom case. Left:  $g^{(2)}(\tau)$  vs the delay time  $\tau$ .  $g^{(2)}(\tau=0)=0.19\pm0.02$  for  $90S_{1/2}$ . Right:  $g^{(2)}(\tau=0)$  vs  $n$  of  $nS_{1/2}$  Rydberg states.

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